Two-gap behaviors of the high- $T_{\rm c}$ cuprate superconductors: Universal versus material-dependent properties

T. Yoshida¹, M. Hashimoto¹, K. Tanaka², N. Mannnella², Z. Hussain³, Z.-X. Shen², A. Fujimori¹, M. Kubota⁴, K. Ono⁴, S. Komiya⁵, Y. Ando⁶, H. Eisaki⁷, S. Uchida¹

¹ Department of Physics, University of Tokyo, Bunkyo-ku, Tokyo 113-0033, Japan

² Department of Applied Physics and Stanford Synchrotron Radiation Laboratory, Stanford University, Stanford, CA94305

Advanced Light Source, Lawrence Berkeley National Lab, Berkeley, CA 94720, USA

⁴Photon Factory, Institute of Materials Structure Science, KEK, Tsukuba, Ibaraki 305-0801, Japan ⁵Central Research Institute of Electric Power Industry, Komae, Tokyo 201-8511, Japan

⁶The Institute of Scientific and IndustrialResearch, Osaka University, Ibaraki, Osaka

567-0047, Japan

⁷National Institute of Advanced Industrial Science and Technology, Tsukuba 305-8568, Japan

One of the central issues in the research of the high- T_c cuprates superconductors is whether the pseudogap is a distinct phenomenon from superconuctivity or a gap due to local pairing or incoherent superconducting fluctuations above T_c . In the former scenario, a possible origin of the pseudogap is preformed Cooper pairs lacking phase coherence. In the latter scenario, the pseudogap is due to a competing order such as spin density wave, charge density wave, d-density wave, etc. It has been well known that the pseudogap in the anti-nodal $(\pi, 0)$ region increases with underdoping as observed by angle-resolved photoemission spectroscopy (ARPES). However, the energy gap measured by other experimental methods such as Andreev reflection, directly associated which is more with superconductivity, exhibits opposite trend, that is, the gap decreases with underdoping, suggesting a different origin of the superconducting gap from the antinodal gap.

A recent ARPES study of Bi₂Sr₂Ca_{1-x}Y_xCu₂O₈ (Bi2212) has revealed the presence of two distinct energy gaps in different regions of momentum space [1,2]. One is the antinodal region as mentioned above, and increases with underdoping. The other opens in the near-nodal region showing a coherent peak, and does not increase with underdoping. On the other hand, attempts have been made to understand the pseudogap within a single d-wave energy gap [3]. In such a single gap picture, preformed Cooper pairs are the most likely origin of the pseudogap.

Since the doping and temperature dependences of the energy gap would reveal the entangled two-gap behavior, we have investigated the energy gap of lightly- to optimally-doped LSCO by ARPES as a function of doping and temperature. In the present work, the momentum dependence of the gap clearly exhibits two-gap behavior as in the case of heavily underdoped Bi2212: the pseudogap Δ^* in the antinodal region and the d-wave like gap Δ_0 around

the node (Fig. 1). The doping dependence of the obtained parameter Δ_0 qualitatively explains the reduction of the $T_{\rm c}$ with underdoping. Furthermore, from comparison of the present results with those on Bi2212 and other cuprates with higher T_c 's, we have found that the magnitude of the Δ^* and the pseudogap temperature T^* is not appreciably material-dependent, suggesting that the pseudogap is properties of a single CuO₂ plane. On the other hand, the magnitude of the Δ_0 , which is proportional to the superconducting gap, is strongly material-dependent (CuO2 layer number dependent) like T_c, suggesting that they are influenced by the effect of neighboring CuO₂ planes and block layers.

References

[1] K. Tanaka et al., Science **314**, 1910 (2006). [2] W. S. Lee et al., Nature 450, 81 (2007). [3] A. Kanigel et al., Nature Physics 2, 447 (2006)

> 40 7 K X=0 07 2 C 30 - N A_{LEM} (meV) 0.4 0.6 0.2 |cos(k,a)-cos(k,a)| / 2 La_{2-x}Sr_xCuO₄ x=0.03 T= 20K 10 x=0.07 x=0.15 Δ₀ 0 0.2 0.4 0.6 0.8 1.0 |cos(k,a)-cos(k,a)| / 2

Fig. 1. Angular dependence of the gap of LSCO observed by ARPES. The definition of Δ^* and Δ_0 is shown for x=0.15.